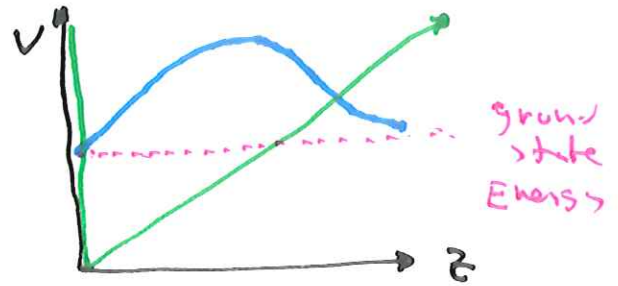


Example: Rayleigh-Ritz - "falling" -  $V(z) = \begin{cases} \infty & z < 0 \\ mgz & z \geq 0 \end{cases}$

$$\psi = z e^{-az}$$

$$\psi' = (1 - az) e^{-az}$$



Dimensionless coordinates:

$$\left( -\frac{\hbar^2}{2m} \right) \partial_z^2 \psi + (mgz) \psi = E \psi$$

A units  $E \cdot L^2$       B units  $\frac{E}{L}$

make length unit  $l = (A/B)^{1/3}$   
 energy unit  $e = (AB^2)^{1/3}$

$$z = z' l$$

$$E = E' e$$

$$\left( -\frac{\hbar^2}{2m l^2} \partial_{z'}^2 + (mg l z') \right) \psi = E' e \psi$$

$$\frac{A}{(A/B)^{2/3}} = (AB^2)^{1/3} = e$$

use:  $\int_0^\infty x^n e^{-ax} dx = \frac{n!}{a^{n+1}}$

$$(-\partial_{z'}^2 + z') \psi = E' \psi$$

$$KE = \langle \psi' | \psi' \rangle = \int_0^\infty (1 - az)^2 e^{-2az} dz$$

$\frac{0!}{2a}$        $\frac{2a \cdot 1!}{(2a)^2}$        $\frac{a^2 \cdot 2!}{(2a)^3}$

$$= \frac{1}{4a}$$

$$PE = \langle \psi | z | \psi \rangle = \int_0^\infty z^3 e^{-2az} dz = \frac{3!}{(2a)^4} = \frac{3}{8a^4}$$

$$N = \langle \psi | \psi \rangle = \int_0^\infty z^2 e^{-2az} dz = \frac{2!}{(2a)^3} = \frac{1}{4a^3}$$

$$\frac{KE + PE}{N} = a^2 + \frac{3}{2a} = 2.476$$

(Note exact answer is 2.338)

$$\frac{d}{da} \left[ \frac{KE + PE}{N} \right] = 0 = 2a - \frac{3}{2a^2} \Rightarrow a^3 = \frac{3}{4}$$

WKB: (Wentzel, Kramers, Brillouin)  $-\psi'' = k^2(x)\psi$

Remark:  $\hbar k(x) = \sqrt{2m(E-V)} = p$   
 $kE = \frac{p^2}{2m}$

$\frac{2m}{\hbar^2} (E - V(x))$

Remark 2: in classically disallowed  $V > E \Rightarrow k(x) = iK$

For an oscillating  $\psi$  try  $\psi = e^{iS(x)}$  where expect  $K$  WOLOG

$S(x_0 + \Delta x) = S(x_0) + S' \Delta x = \text{const} + 2\pi \frac{\Delta x}{\lambda}$

Here we have a "changing wave length" so  $\lambda(x)$

$\psi = e^{iS}$

$\psi' = iS' e^{iS}$

$\psi'' = [(iS')^2 + iS''] e^{iS} = -k^2(x)\psi$

assume "small" so  $S' = \pm k \Rightarrow S = \int^x k(x) dx$

WOLOG:  $S = \int^x k(x) dx + \epsilon$

$S' = k(x) + \epsilon'$

$S'' = k' + \epsilon''$

$[iS'' - S'^2] = [i(k' + \epsilon'') - (k + \epsilon')^2] = -k^2$

$i k' - 2k\epsilon' - \underbrace{\epsilon'^2 + i\epsilon''}_{\text{neglect}} = 0$

$\Rightarrow \epsilon' = \frac{i k'}{2k} \Rightarrow \epsilon = \frac{i}{2} \ln(k) + \text{const}$

$= \frac{i}{2} \ln(\sqrt{k}) + \text{const}$

$\Rightarrow \psi = e^{iS} = e^{i \int k dx} e^{-\ln(\sqrt{k})} = \frac{1}{\sqrt{k}} e^{i \int k dx}$

amplitude large as  $k \rightarrow 0$  (turning pt)

$k(x) = \frac{2\pi}{\lambda(x)}$

Remark: if in classically disallowed  $k = iK \Rightarrow$

$\psi = \frac{1}{\sqrt{k}} e^{\pm \int K dx}$

exponential growth/decay

Qualitative wavefunction - Big amplitude where particle slow.  
 When slow  $\lambda = \frac{h}{p}$  Big  $\lambda$ . Expo decay into disallowed region

Basic Idea (Bohr)  
 Fit an integer number of half wavelengths between turning points

$$\Rightarrow \int_{t_p}^{t_p'} k(x) dx = n\pi$$

$$k dx = 2\pi \frac{dx}{\lambda}$$

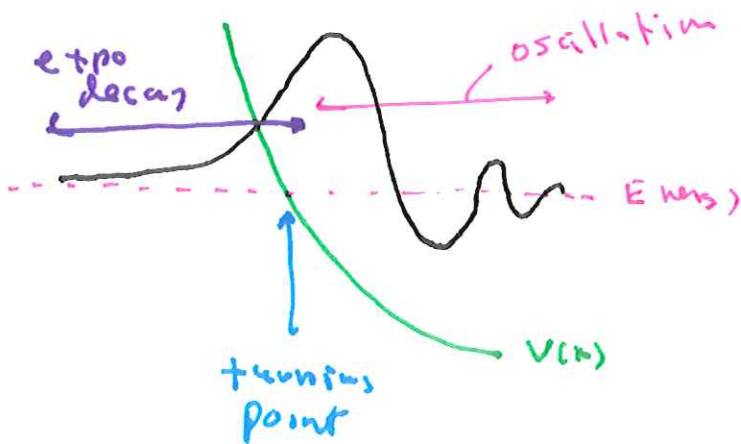
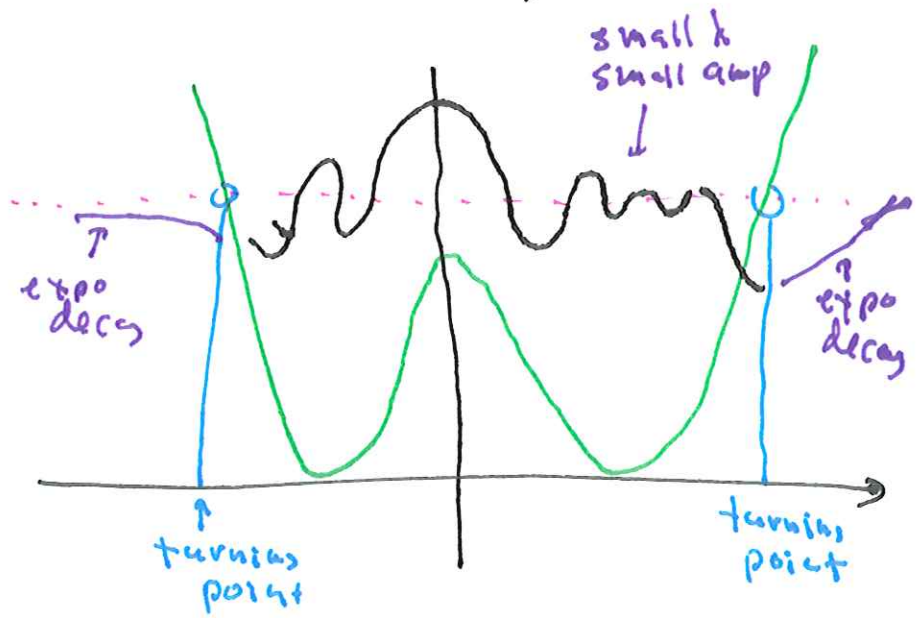
Fraction of wavelength

$$\Rightarrow \int_{t_p}^{t_p'} k(x) dx = 2\pi \text{ (total fraction of wavelength)}$$

$$\text{make} = \frac{\text{integer}}{2}$$

Update Idea - wave function not actually zero at turning point [ie doesn't start as zero has already risen from expo decay]

unless  $V = \infty$



Turns out at turning point its already at  $45^\circ = \frac{1}{8}$  wave  
 → with turning points at both ends  $\frac{1}{4}$  wave missing

⇒ seek integer # half waves minus  $\frac{1}{4}$  wave

$$\Rightarrow \int k(x) dx = \pi \left( n - \frac{1}{2} \right)$$

# half waves

correction:  $\frac{1}{4}$  wave

