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Perturbation Theory: $H\Psi = E\Psi$ where $H = H_0 + V$ where $H_0\Psi_0 = E_0\Psi_0$ (known solution) and V is the perturbation. Parameter: $0 \rightarrow 1$

Find Taylor series expansion for E

$$E = E_0 + \lambda E_1 + \lambda^2 E_2 + \dots$$

E_0 is some unperturbed energy E_n with eigenfunction Ψ_n

$$\sum_{m \neq n} \frac{|\langle m | V | n \rangle|^2}{E_n - E_m}$$

Problem: if H has degeneracy then this term may be zero... if E_2 is infinite E_1 is meaningless

Solution requires $\langle n | V | n \rangle = 0$ if $E_n - E_m = 0$ but that will not generally be true - we must make it true by redefining the states (wavefunctions) that are degenerate.

Words: degenerate subspace: the space spanned by Ψ_n and its degenerate wavefunctions.

In the H atom all the states with the same value of n are degenerate as E does not depend on l or m

eg $n=3$: 3s, 3p (for $m=1, 0, -1$), 3d (for $m=2, 1, 0, -1, -2$)

These $1+3+5=9$ states are degenerate

If we want to determine how these states are changed in the presence of electric field we need to find a new basis (made up of linear combos of the above states) that has the property

$$\langle a | V | B \rangle = 0$$

these are our new basis states in contrast to the original set [2s, 3p ($m=1, 0, -1$), etc] which we will label i, j, k etc

Summary: A subset of the original unperturbed wavefunction $\Psi_n = |n\rangle$ are degenerate (i.e. same E) which generally results in $E_2 = \infty$ which means E_1 is nonsense. The solution is to find a new set of states by a linear combo of the degenerate $|n\rangle$. These new states will still be degenerate but we aim to make them so that $\langle \alpha | V | \beta \rangle = 0$ which keeps E_2 finite.

one way to do this is by making the states $|a\rangle$ from an eigenvector of the matrix $\langle i | V | j \rangle = V_{ij}$ if $|a\rangle = \sum c_i |i\rangle$ where $\begin{bmatrix} V_{ij} \end{bmatrix} \begin{bmatrix} c \end{bmatrix} = E_a \begin{bmatrix} c \end{bmatrix}$

and $|\beta\rangle = \sum d_i |i\rangle$ where $\begin{bmatrix} V_{ij} \end{bmatrix} \begin{bmatrix} d \end{bmatrix} = E_b \begin{bmatrix} d \end{bmatrix}$

$$\text{Then } \langle \alpha | V | \beta \rangle = \sum_{i,j} c_i^* \langle i | V | j \rangle d_j = E_b \sum_i c_i^* d_i = 0$$

if $E_a = E_b$
 no help } as eigenvectors with different eigenvalues are orthogonal

Note: under these conditions first order

$$E_1 = \langle \alpha | V | \alpha \rangle = E_a \sum c_i^* c_i = E_a$$

So first order correct are these eigen values

actually "Gram-Schmidt" says with a bit of work this case can also be made orthogonal - see p 440

Another way to achieve $\langle \psi | V | \psi \rangle = 0$ is to find a third operator Q such that $[H, Q] = 0$

if $[H, Q] = 0$ we can select ψ_n such that $[V, Q] = 0$

to be eigen functions of Q & H : $H \psi_n = E_n \psi_n$
 $Q \psi_n = q_n \psi_n$

then: $0 = \langle i | [V, Q] | j \rangle = \langle i | V q_j | j \rangle - \langle i | q_i V | j \rangle$
 $= (q_j - q_i) \langle i | V | j \rangle$

So if $q_j \neq q_i$ $\langle i | V | j \rangle = 0$

Example: Stark effect ($V = z$) for the 9 degenerate states of $n=3$ H-atom 9
12-27

$|l, m\rangle$: 1 0 0 1 1 1 1 1 0 1 1 -1 1 2 2 1 2 1 1 2 0 1 2 -1
1 2 3 4 5 6 7 8

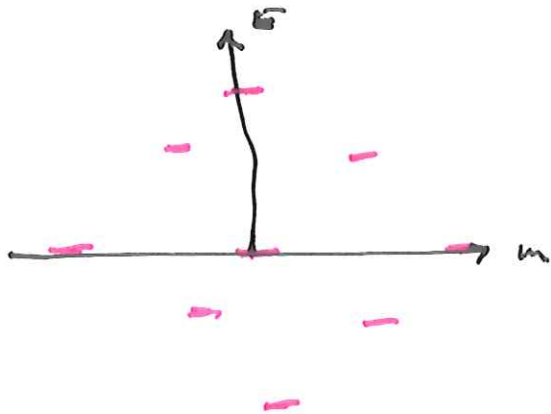
Using Mathematica we find the 9×9 matrix $\langle i | z | j \rangle$ and its eigen values / vectors

values: $-9, 9, -9/2, 9/2, 9/2, 9/2, 0, 0, 0$

- vectors:
- $(\sqrt{2}, 0, \sqrt{3}, 0, 0, 0, 1, 0, 0) \leftarrow m=0$
 - $(\sqrt{2}, 0, -\sqrt{3}, 0, 0, 0, 1, 0, 0) \leftarrow m=0$
 - $(0, 0, 0, 1, 0, 0, 0, 1, 0) \leftarrow m=0$
 - $(0, 1, 0, 0, 0, 0, 1, 0, 0) \leftarrow m=1$
 - $(0, 0, 0, -1, 0, 0, 0, 1, 0) \leftarrow m=-1$
 - $(0, -1, 0, 0, 0, 0, 1, 0, 0) \leftarrow m=1$
 - $(0, 0, 0, 0, 0, 0, 0, 0, 1) \leftarrow m=2$
 - $(\frac{1}{\sqrt{2}}, 0, 0, 0, 0, 0, 1, 0, 0) \leftarrow m=0$
 - $(0, 0, 0, 0, 0, 0, 1, 0, 0) \leftarrow m=2$

putting things together:

$m=0$	$-q, q, 0$
$m=1$	$-q/2, +q/2$
$m=2$	$-q/2, +q/2$
$m=-1$	0
$m=-2$	0



Note: There are for example 3 states with eigenvalue = 0
 we are not guaranteed these satisfy $\langle \alpha | V | \beta \rangle = 0$
 BUT these 3 states have different m values
 \ddagger $[L_z, H] = 0$ $[L_z, V] = 0$ so the different
 $"q"$ values guarantees $\langle \alpha | V | \beta \rangle = 0$

Note: you can just do the dot products with the
 3 $m=0$ vectors and see they are orthogonal
 (but not normalized yet)