

math vector space over a field of scalars

→ addition that is commutative, associative,  $\exists$  inverses

→ scalar multiplication that is commutative, associative

Note: physics "vector" carries many additional properties + math "vector"

basis vectors  $|e_i\rangle$  that span space; dimension

$$|\vec{v}\rangle = \sum q_i |e_i\rangle \quad |\vec{v}\rangle \leftrightarrow \begin{pmatrix} q_1 \\ \vdots \\ q_n \end{pmatrix} \text{ "ket"}$$

Additional property: inner product ("dot product")

$$\langle \vec{w} | \vec{v} \rangle = (\langle \vec{v} | \vec{w} \rangle)^* ; \langle \vec{v} | \vec{v} \rangle \geq 0 \quad ; \quad = 0 \text{ iff } \vec{v} = \vec{0}$$

satisfies linear combinations

orthonormal basis:  $\langle e_i | e_j \rangle = \delta_{ij}$

$$\langle \vec{w} | = (b_1^*, \dots, b_n^*) ;$$

↑ "bra"

$$\langle \vec{w} | \vec{v} \rangle = \underbrace{(b_1^*, \dots, b_n^*)}_{\text{bracket}} \underbrace{\begin{pmatrix} q_1 \\ \vdots \\ q_n \end{pmatrix}}_{\text{ket}} = \sum b_i^* q_i$$

operator T:

$$T |e_i\rangle = \sum_j T_{ji} |e_j\rangle$$

we use this exact definition so things work out in future

composition (product) of 2 operators [like  $g(f(x))$ ]

$$\begin{aligned} ST |e_i\rangle &= \sum_j T_{ji} S |e_j\rangle = \sum_j T_{ji} \sum_k S_{kj} |e_k\rangle \\ &= \sum_k \left( \sum_j S_{kj} T_{ji} \right) |e_k\rangle \end{aligned}$$

so matrix that represents (ST) is just the matrix product of the matrices that represent S & T

Note: transpose of matrix on vector  $\sim$  in Griffiths  
hermitian conjugate (adjoint) † (dagger)

see  $\langle \vec{w} | \vec{v} \rangle = (\langle \vec{w} |)^{\dagger} | \vec{v} \rangle$

$$\langle \vec{w} | T \vec{v} \rangle = (T^{\dagger} \langle \vec{w} |)^{\dagger} | \vec{v} \rangle = \langle T^{\dagger} \vec{w} | \vec{v} \rangle$$

cuz:  $(AB) = \tilde{B} \tilde{A} \quad ; \quad T^{\dagger\dagger} = T$

define: Hermitian:  $T^{\dagger} = T$  ; Unitary  $T^{\dagger} = T^{-1}$

FYI: Symmetric  $\tilde{T} = T$  ; orthogonal  $\tilde{T} = T^{-1}$

Eigen vectors & Eigen values:  $T\vec{v} = \lambda\vec{v}$  "characteristic"

[example: the "principal axes of rotation" were axes for which the angular momentum  $\vec{L}$  was in same direction as angular velocity  $\vec{\omega}$ ; i.e.,  $\vec{L} = \vec{I}\vec{\omega} = \lambda\vec{\omega}$  ]

if  $T\vec{v} = \lambda\vec{v}$  then  $(T - \lambda\vec{1})\vec{v} = 0$   
↑ unit matrix

so  $T - \lambda\vec{1}$  has a null vector & hence has NO inverse.

$$\text{so } \det(T - \lambda\vec{1}) = 0$$

For Hermitian matrices: the eigen vectors are orthogonal<sup>②</sup>; span space; eigen values are real<sup>①</sup>

$$\textcircled{1}: \langle \vec{v} | T\vec{v} \rangle = \lambda \langle \vec{v} | \vec{v} \rangle = \langle T\vec{v} | \vec{v} \rangle = \lambda^* \langle \vec{v} | \vec{v} \rangle$$

$$\textcircled{2} \left. \begin{array}{l} T\vec{v} = \lambda\vec{v} \\ T\vec{w} = \mu\vec{w} \end{array} \right\} \begin{array}{l} \langle \vec{w} | T\vec{v} \rangle = \lambda \langle \vec{w} | \vec{v} \rangle \\ \langle T\vec{w} | \vec{v} \rangle = \mu \langle \vec{w} | \vec{v} \rangle \end{array}$$

$$\text{so, } \underbrace{(\lambda - \mu)}_{\neq 0} \langle \vec{w} | \vec{v} \rangle = 0$$

Case of "degenerate eigenvalues" via

"Gram-Schmidt" process

If we use eigen vectors as basis, matrix represent

T is simple: diagonal

$$T|e_i\rangle = \lambda_i|e_i\rangle \Rightarrow T =$$

$$\begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix}$$

"diagonalization"

How does the matrix that represents an operator change if we select a different basis?

[eg How would the moment of inertia tensor change if we selected a rotated coordinate system?]

Begin by writing the original basis vectors as a linear combination of the new basis (F) - we pick a particular naming for these expansion coeffs:

$$|e_i\rangle = \sum_j S_{ji} |f_j\rangle \longrightarrow \sum_i (S^{-1})_{ik} |e_i\rangle = |f_k\rangle$$

Matrix representation T in new basis:  $(T^F)_{ij} = \langle f_i | T | f_j \rangle$

$$\begin{aligned} T | f_j \rangle &= T \sum_k (S^{-1})_{kj} |e_k\rangle = \sum_k (S^{-1})_{kj} T |e_k\rangle \\ &= \sum_{k,l} (S^{-1})_{kj} T_{lk} |e_l\rangle \\ &= \sum_{k,l,n} (S^{-1})_{kj} T_{lk} S_{nl} |f_n\rangle \end{aligned}$$

$$\begin{aligned} \text{So } \langle f_i | T | f_j \rangle &= \sum_{k,l} (S^{-1})_{kj} T_{lk} S_{il} \\ &= \sum_{k,l} \underbrace{S_{il} T_{lk} (S^{-1})_{kj}}_{\text{matrix } S T S^{-1}} \end{aligned}$$

Notes: if this is a coordinate rotation S will be an orthogonal matrix

if this is a change of orthonormal basis S will be unitary matrix

The determinant & trace of a matrix are invariant under a similarity transformation

Eigenvalues also invariant.